

**Physics 4062/5062 – Lecture Eight – Strong Field Solution**

$$i\dot{c}_1 = \Omega \cos(\omega t) e^{i\omega_0 t} c_2 \quad (207)$$

$$i\dot{c}_2 = \Omega^* \cos(\omega t) e^{i\omega_0 t} c_1$$

- Writing the cosine in exponential notation and using the rotating wave approximation, we obtain

$$i\dot{c}_1 = c_2 \frac{\Omega}{2} \exp(i(\omega - \omega_0)t) \quad (213)$$

$$i\dot{c}_2 = c_1 \frac{\Omega^*}{2} \exp(-i(\omega - \omega_0)t)$$

- combine equations in (213)

$$\frac{d^2 c_2}{dt^2} + i(\omega - \omega_0) \left( \frac{dc_2}{dt} \right) + \left| \frac{\Omega}{2} \right|^2 c_2 = 0 \quad (214)$$

- Use initial conditions

$$c_1(0) = 1$$

$$c_2(0) = 0$$

- To obtain:

$$|c_2(t)|^2 = \left( \frac{\Omega^2}{\Omega'^2} \right) \sin^2 \left( \frac{\Omega' t}{2} \right) \quad (215)$$

$$\Omega'^2 = \Omega^2 + \Delta^2 \quad (216) \quad (\text{Generalized Rabi Frequency})$$

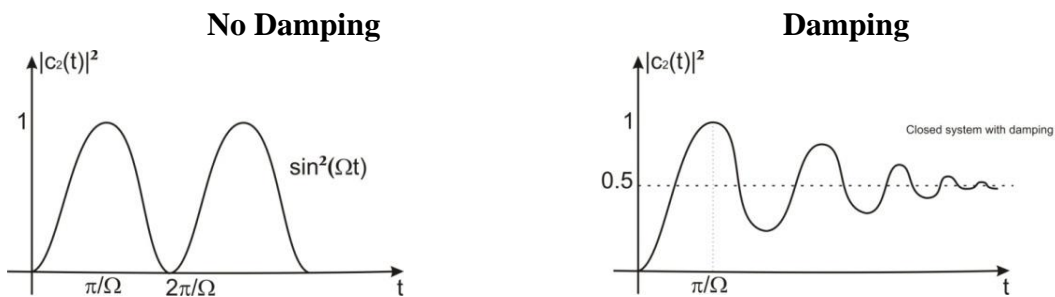
- At Resonance

$$\omega = \omega_0$$

$$\Omega' = \Omega$$

$$|c_2(t)|^2 = \sin^2 \left( \frac{\Omega t}{2} \right) \quad (217)$$

Population Oscillates between 2 levels at characteristic flopping frequency ( $\Omega$ )



- For  $\Omega t = \pi \Rightarrow |c_2(t)|^2 = 1$  population inverted
- Behavior is different from two level system described by rate equations
- In the rate equation approach, populations of levels 1 and 2 become equal with increase in excitation rate and there is no inversion

### Bloch Vector Model

Define Dipole Moment in terms of expectation value

$$p = \int \psi^*(t)(qx)\psi(t)d^3r \quad (218)$$

Use  $\psi(r, t) = c_1\psi_1(r)e^{-\frac{iE_1t}{\hbar}} + c_2\psi_2(r)e^{-\frac{iE_2t}{\hbar}}$  in equation 218 to obtain

$$x(t) = c_2^*c_1x_{21}e^{i\omega_0t} + c_1^*c_2x_{12}e^{-i\omega_0t} \quad (219)$$

Here,  $x_{12} = \langle 1|\hat{x}|2\rangle = x_{12}^*$

$$\omega_0 = \omega_1 - \omega_2$$

$$x_{11} = x_{22} = 0$$

**Goal:** Find atomic populations and components of dipole moment induced by external field

Introduce Density Matrix for ensemble of atoms

$$|\psi\rangle\langle\psi| = \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} (c_1^* \quad c_2^*) = \begin{pmatrix} |c_1|^2 & c_1c_2^* \\ c_2c_1^* & |c_2|^2 \end{pmatrix} = \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} \quad (220)$$

$\rho_{11}$  and  $\rho_{22}$  are the populations

$\rho_{12}$  and  $\rho_{21}$  are coherences associated with dipole moment  $qx(t)$  where  $x(t)$  is given by equation 219

**Introduce variable change**

$$\tilde{c}_1 = c_1 \exp[-i\Delta t/2]$$

$$\tilde{c}_2 = c_2 \exp[-i\Delta t/2]$$

As a result, populations are unaffected, that is

$$\tilde{\rho}_{11} = \rho_{11}$$

$$\tilde{\rho}_{22} = \rho_{22}$$

The new coherences are given by

$$\tilde{\rho}_{12} = \rho_{12} e^{-i\Delta t} \quad (223)$$

$$\tilde{\rho}_{21} = \rho_{21} e^{i\Delta t}$$

The atomic response becomes  $x(t) = x_{12}[\tilde{\rho}_{12} \exp[i\omega t] + \tilde{\rho}_{21} \exp[-i\omega t]]$  (224)

Define in phase and quadrature components of dipole moment

$$\begin{aligned} \mathcal{U} &= \tilde{\rho}_{12} + \tilde{\rho}_{21} \\ \mathcal{V} &= -i(\tilde{\rho}_{12} - \tilde{\rho}_{21}) \end{aligned} \quad (225)$$

As a result, equation 224 becomes  $x(t) = x_{12}(\mathcal{U}(t) \cos(\omega t) - \mathcal{V}(t) \sin(\omega t))$

Note that  $\mathcal{U}$  and  $\mathcal{V}$  are components of dipole moment in frame rotating at  $\omega$

To find  $\tilde{\rho}_{12}$ ,  $\tilde{\rho}_{21}$  and  $\rho_{22}$

Rewrite 213 in the new notation

$$\begin{aligned} \frac{dc_1}{dt} &= c_2 \exp[i\Delta t] \frac{\Omega}{2} \\ \frac{dc_2}{dt} &= c_1 \exp[-i\Delta t] \frac{\Omega}{2} \end{aligned} \quad (226 \text{ A}) \text{ and } (226 \text{ B})$$

From 221,

$$\frac{d\tilde{c}_1}{dt} = \frac{dc_1}{dt} \exp[-i\Delta \frac{t}{2}] - i \frac{\Delta}{2} c_1 \exp[-i\Delta \frac{t}{2}]$$

Multiply by  $i$  and use equations 226 A, 221 and 222

Multiply by  $i$  and use equations 226 B, 221 and 222 to obtain coupled equations

$$\begin{aligned} \frac{d\tilde{c}_1}{dt} &= \frac{1}{2}[\Delta\tilde{c}_1 + \Omega\tilde{c}_2] \\ \frac{d\tilde{c}_2}{dt} &= \frac{1}{2}[\Omega\tilde{c}_1 - \Delta\tilde{c}_2] \end{aligned}$$

Find  $\frac{d\tilde{\rho}_{12}}{dt} = \tilde{c}_1 \frac{d\tilde{c}_2^*}{dt} + \frac{d\tilde{c}_1}{dt} \tilde{c}_2^*$ , and  $\frac{d\tilde{\rho}_{22}}{dt}$  using the normalization condition  $\rho_{11} + \rho_{22} = 1$

It can be shown that the following coupled equations are obtained

$$\begin{aligned}\dot{u} &= \Delta\mathcal{V} \\ \dot{v} &= \Delta\mathcal{U} + \Omega(\rho_{11} - \rho_{22}) \quad (229) \\ \dot{\rho}_{22} &= \frac{\Omega\mathcal{V}}{2}\end{aligned}$$

**Define:**

$$\mathcal{W} = \rho_{11} - \rho_{22} \quad (230)$$

As a result the following optical Bloch equations are obtained.

$$\begin{aligned}\dot{u} &= \Delta\mathcal{V} \\ \dot{v} &= -\Delta\mathcal{U} + \Omega\mathcal{W} \quad (231) \\ \dot{\mathcal{W}} &= -\Omega\mathcal{V}\end{aligned}$$

Define Bloch Vector that describes a point on the Bloch sphere of radius  $|\mathcal{U}|^2 + |\mathcal{V}|^2 + |\mathcal{W}|^2 = 1$

$$\vec{R} = u\hat{1} + v\hat{2} + \mathcal{W}\hat{3}$$

Define interaction with field

$$\vec{\Omega}' = \Omega\hat{1} + \Delta\hat{3}$$

With these definitions, equation 231 is reduced to torque equation

$$\dot{\vec{R}} = \vec{R} \times \vec{\Omega}'$$

Model describes precession of Bloch vector around Omega Prime (vector symbol).

**Time dependent Components of Bloch vector give populations and coherences**

